

Dissipative cooling of spin chains
by a bath of dipolar particles

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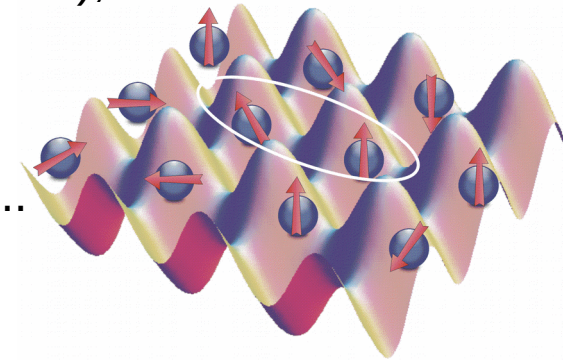
Magnetism with cold atoms

Various magnetic models implemented in cold atoms

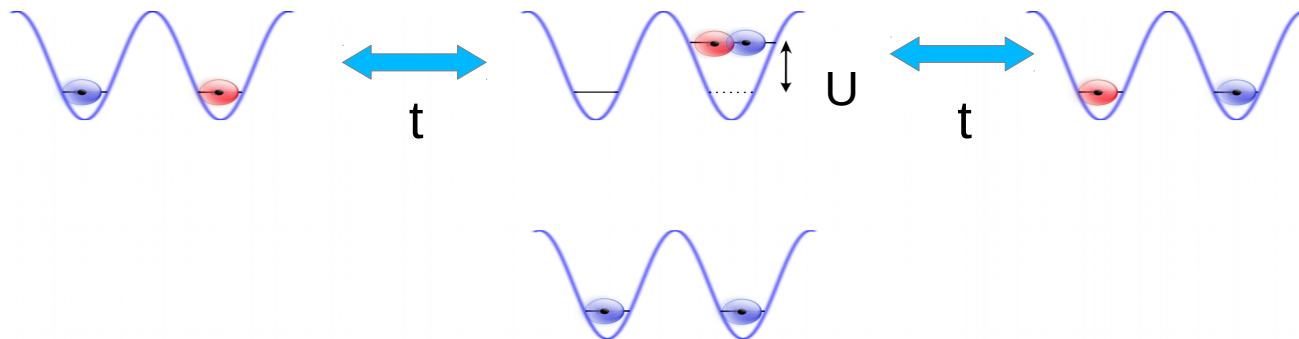
Broad panel of physical questions : frustration (tunable geometries), large spin systems, interplay with transport (t-J model), ...

Variety of magnetic interactions using ground state atoms, Rydberg state atoms, molecules, mappings... (spin-dependence, short- or long-range, anisotropy)

→ Heisenberg, Ising, XXZ, and others...



Much studied : antiferromagnetic Heisenberg model from super-exchange in the Mott regime



$$H = - J \sum_{\langle i, j \rangle} \vec{S}_i \cdot \vec{S}_j$$

with $J \approx -4t^2/U$

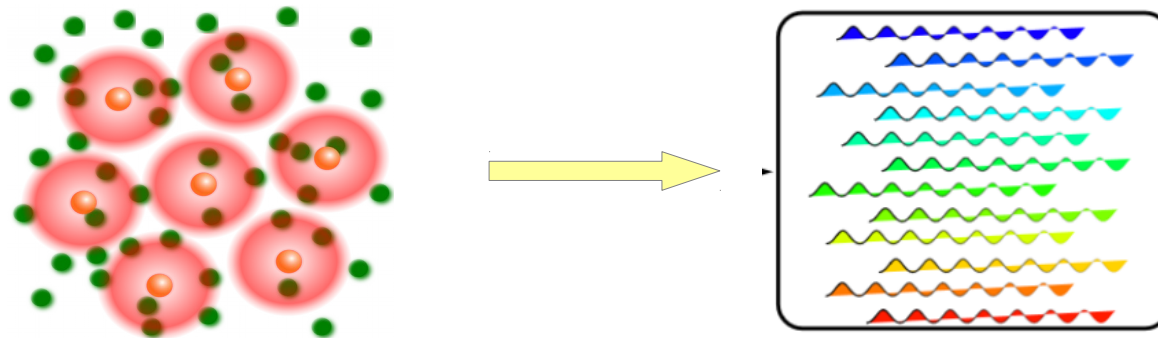
An opportunity for open quantum systems

Many-body quantum systems = resources for quantum simulation and quantum computation

Dissipation ?

Modern research field :
robust entanglement / correlations from engineered dissipation

Benatti et al, PRL 91, 070402 (2003) - Piani, Zoller, Cirac, ...
Rydbergs, ions



Many-body system, e.g. Rydberg gas

Bath continuum (e.g., light)

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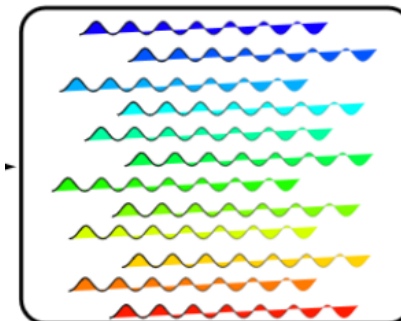
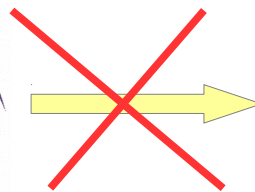
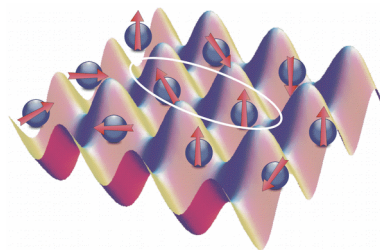
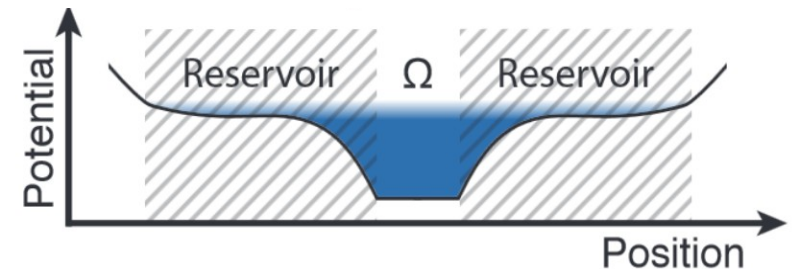
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Rydbergs, ions

Ground state lattice magnetism usually in isolated systems

Strong motivation: the low entropy challenge (Mc Kay and DeMarco, 2011)

so far tackled by **inhomogeneous systems**

Ho 2009, Bernier 2009, Mathy 2012,
Hart 2015, Mazurenko 2017, Kantian 2018 ...



Ground state lattice gas

Bath continuum (e.g., light)

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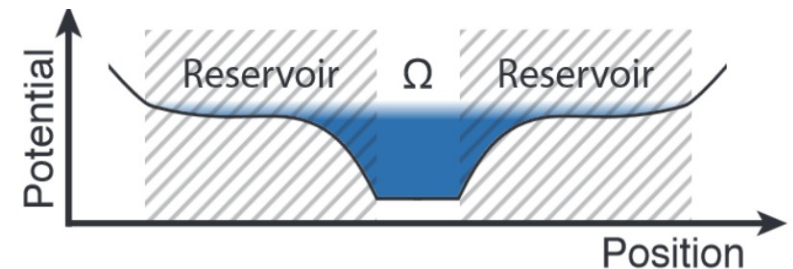
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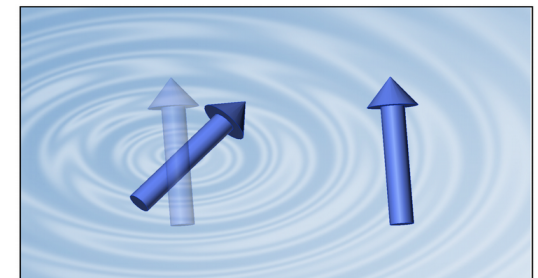
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Engineered spin dissipation proposals : growing literature proposing light as bath

Diehl 2010, Kaczmarczyk 2016, ...
Zoller, Weimer, ...

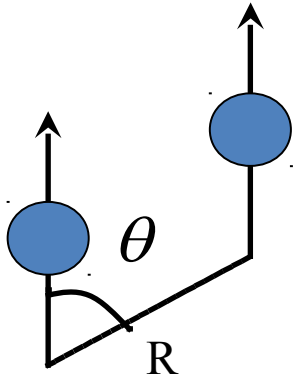
Our discussion:
thermalize the spins with the phonons of an atomic bath
(atomic mixtures)



The tool: dipolar interactions

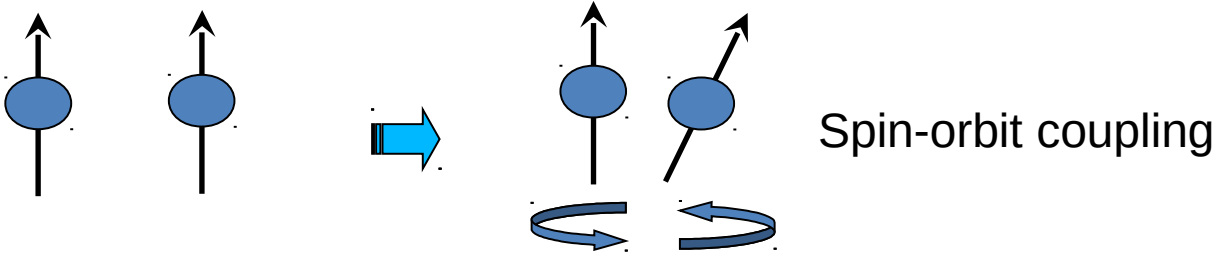
The bath must be able to flip spins

Magnetic dipolar interactions – anisotropic – non spin conserving



$$V_{dd} = \frac{\mu_0}{4\pi} (g_J \mu_B)^2 \frac{S_1 \cdot S_2 - 3(S_1 \cdot u_R)(S_2 \cdot u_R)}{R^3}$$

Dipolar quantum gases:
Pfau, Laburthe-Tolra,
Lev, Ferlaino, Grimm,
Inguscio...



Ising

Exchange

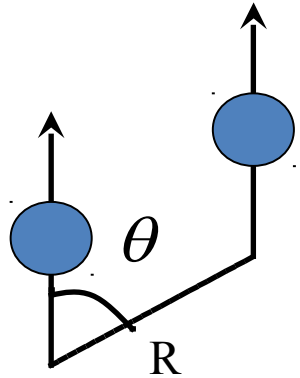
$$S_{1z} S_{2z} + \frac{1}{2} (S_{1+} S_{2-} + S_{1-} S_{2+}) - \frac{3}{4} (2z S_{1z} + r_- S_{1+} + r_+ S_{1-}) (2z S_{2z} + r_- S_{2+} + r_+ S_{2-})$$

Includes non-spin conserving terms

The tool: dipolar interactions

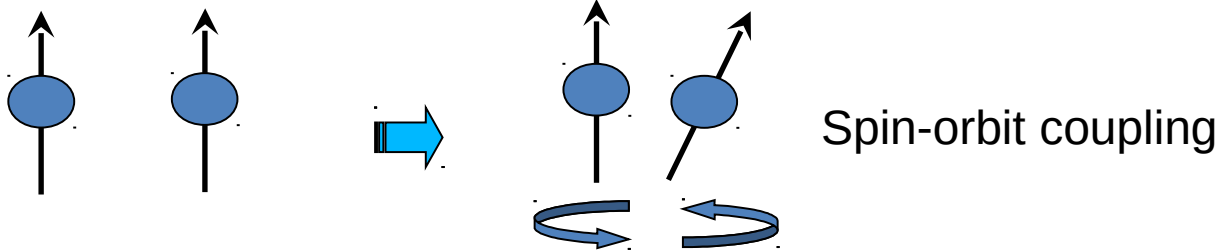
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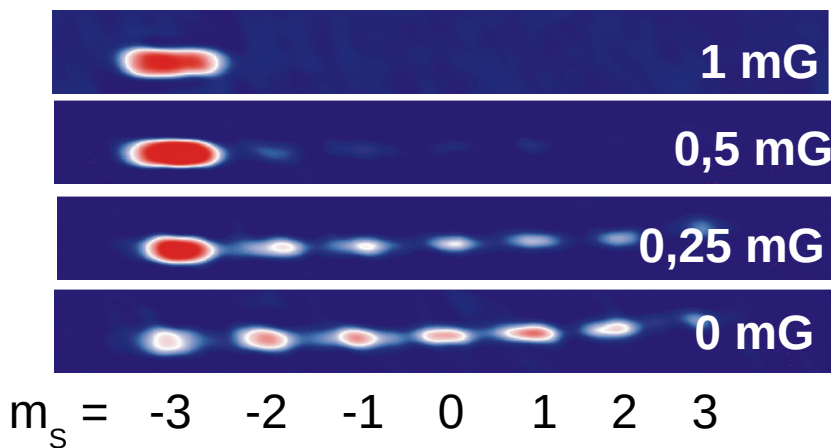
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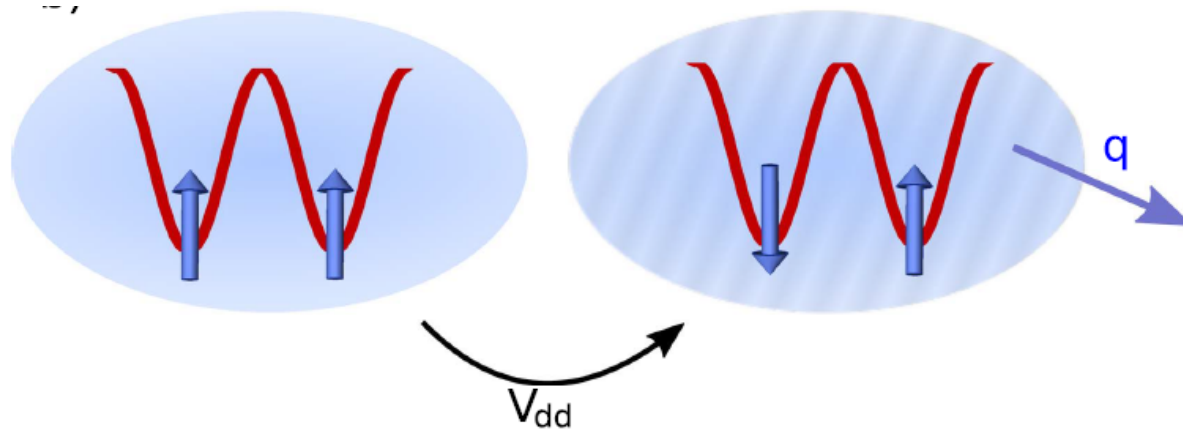
Dipolar relaxation enables true thermalization at free magnetization

Single-species Chromium experiment at LPL (Laburthe-Tolra)



Pasquiou et al, PRL **106**, 255303 (2011)

The gaz always reaches the energetically-favourable spin distribution



This talk:

Dipolar interactions between a spinfull Mott insulator and a dipolar BEC offer true thermalization of the spin degree of freedom of the Mott insulator

- **free magnetization** from a spin-orbit coupling mechanism
- dissipative preparation / protection of highly correlated states

Timescales : compatible with alkali spin chains

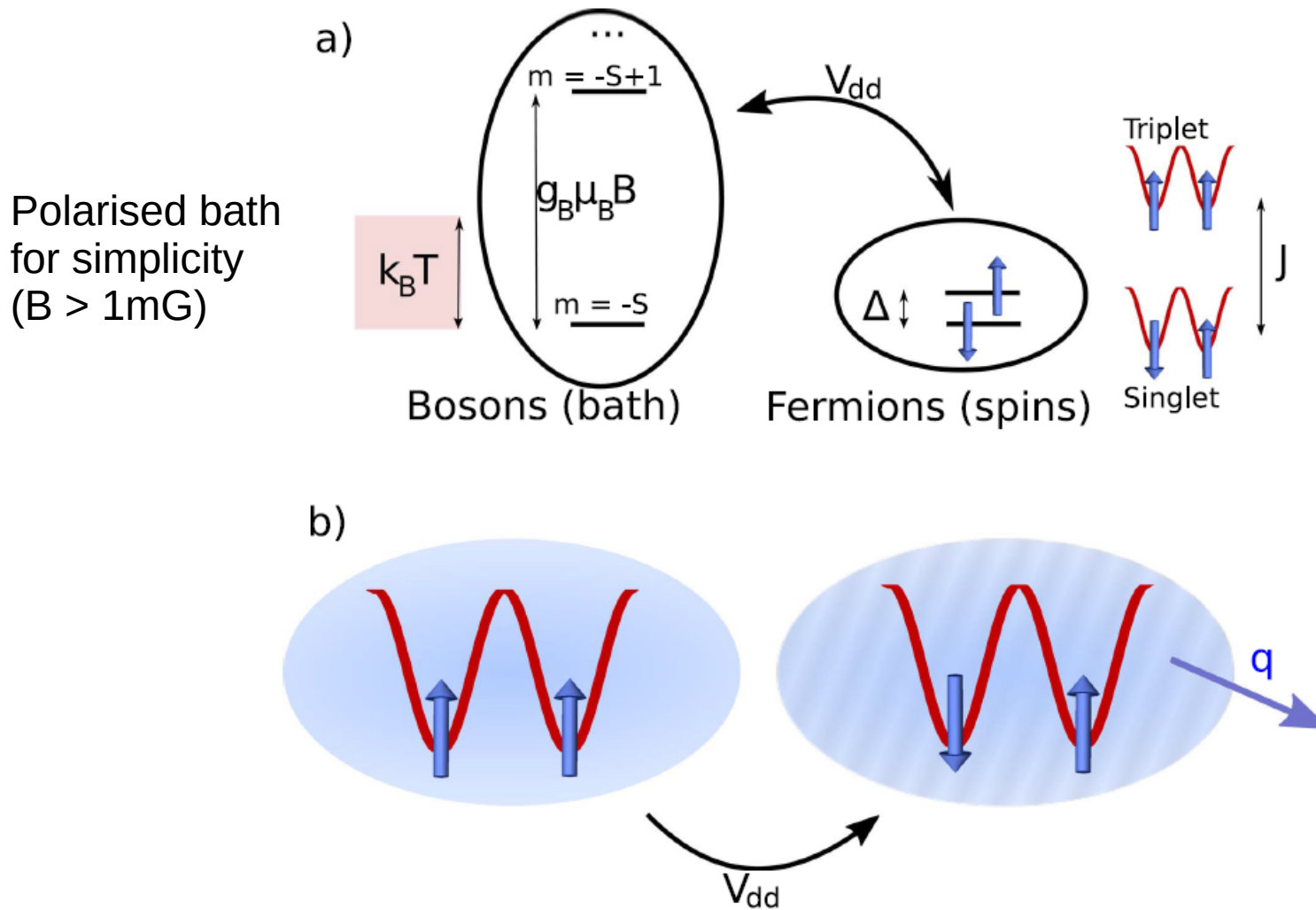
I) Overview of the physics

System overview
A Fermi Golden rule treatment
Anisotropic coupling to the bath

II) Realistic system – numerical calculation

Lattice potential effect on the bath
Convergence to a thermal state
Collective spin dynamics

System overview and simplifying assumptions



AF Heisenberg model;

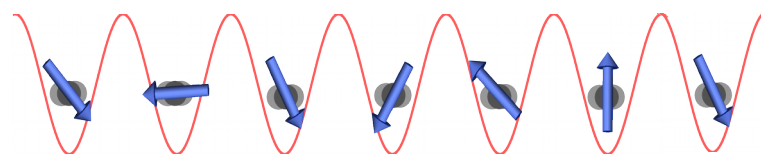
Restricted to two Zeeman states for simplicity

Degenerate

(see NJP appendix, and Gerbier 2006, PRA **73**, 041602)

Bath: Bogoliubov description in the lattice; finite temperature.

Spin chains : finite size 1D chain (up to 7), exactly diagonalized, neglecting any hole/doublon

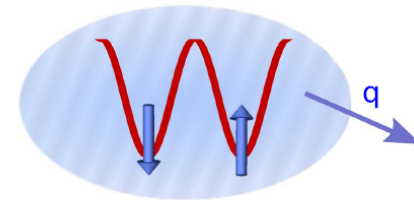


A Fermi Golden rule treatment

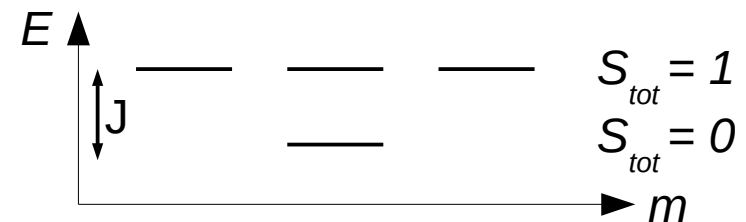
Dissipative evolution evaluated from the Fermi golden rule between collective spin chain eigenstates

$$\Gamma_{i \rightarrow f} = \frac{2\pi}{\hbar} \sum_{|f_{\text{bath}}\rangle} |\langle f_{\text{spin}}; f_{\text{bath}} | H_{\text{int}} | i_{\text{spin}}; i_{\text{bath}} \rangle|^2 \delta(E_{if} + E_{if}^{\text{bath}})$$

$$\frac{dp_i}{dt} = \sum_f (-\Gamma_{i \rightarrow f} p_i + \Gamma_{f \rightarrow i} p_f)$$



Example : 2-atom spin chain, four collective states



Our work: Compute explicitly all these matrix elements, in realistic setting

Detailed calculation in NJP 20, 073037 (2018)

Species: alkali + dipolar. Here, ^{40}K as spin chain, ^{164}Dy as highly dipolar species ($10 \mu_B$)*

Anisotropic coupling to the bath

Radiation diagrams from two spins (double well)

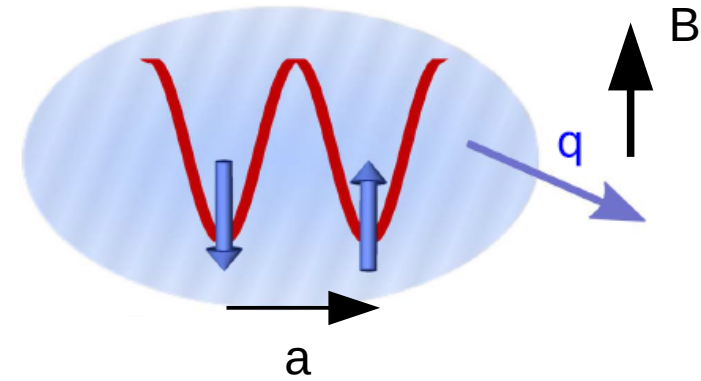
(here without lattice potential for the bath)

Severe anisotropy and collective spin dependence

Example : rate from $S_{\text{tot}} = 1, m = 0$ to $S_{\text{tot}} = 0, m = 0$

Sz component of $|\langle f_{\text{spin}}; f_{\text{bath}} | H_{\text{int}} | i_{\text{spin}}; i_{\text{bath}} \rangle|^2$
(m conserving)

$$|f_{\text{bath}}(\vec{q})\rangle = b^\dagger(\vec{q})|\text{BEC}\rangle \quad \text{with} \quad \epsilon(\vec{q}) = J$$

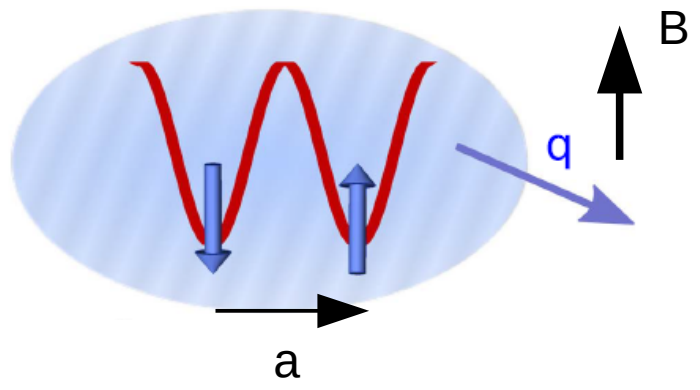


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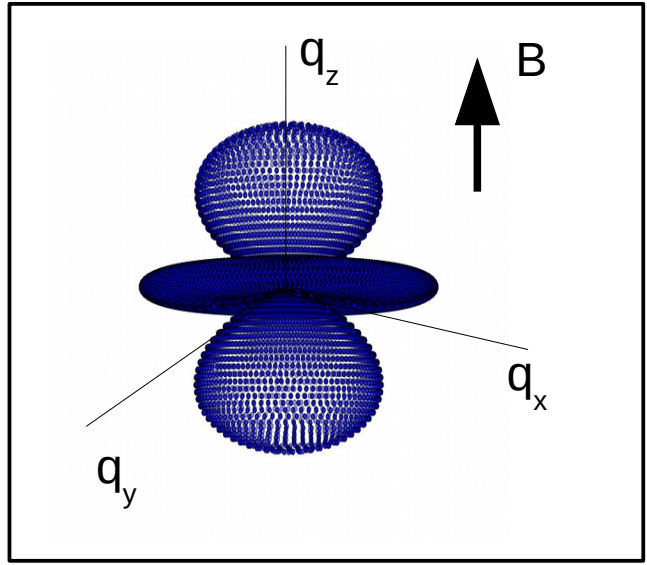


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Single spin : Vdd(q)



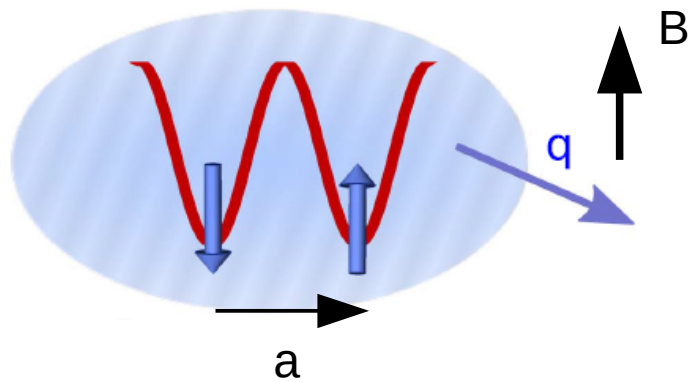
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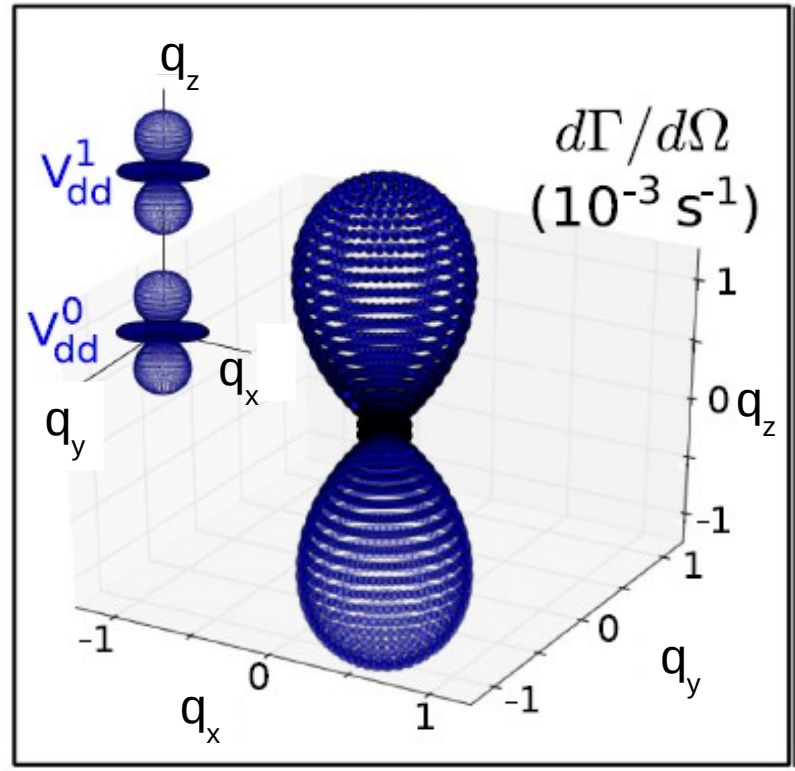
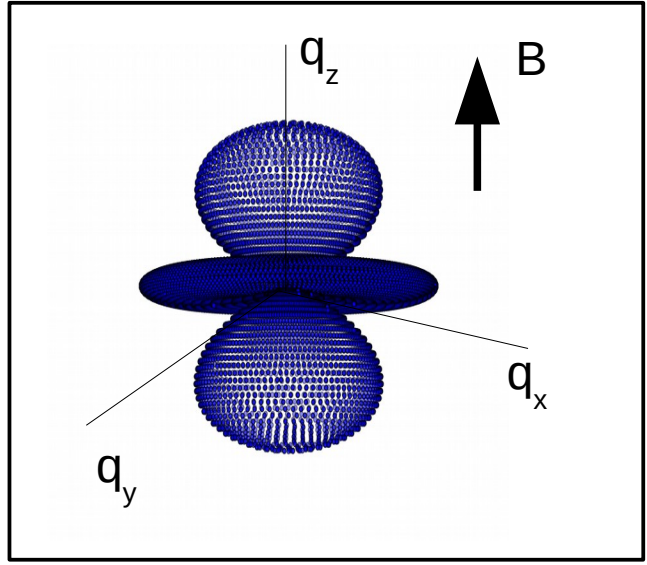


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Two spins $|\langle f | H_{int} | i \rangle|^2$

Single spin : $V_{dd}(\mathbf{q})$



$\vec{q} \cdot \vec{a} = 0$: global energy shift, no effect

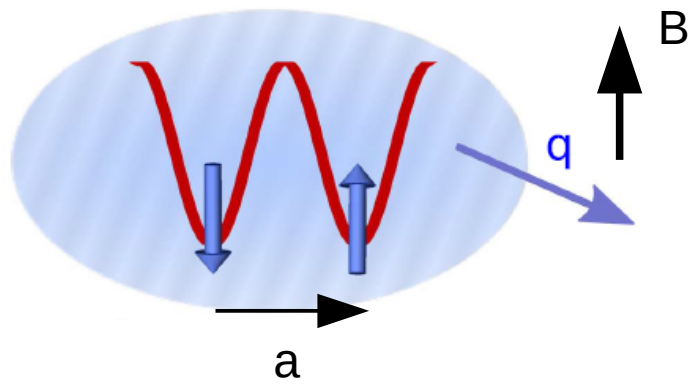
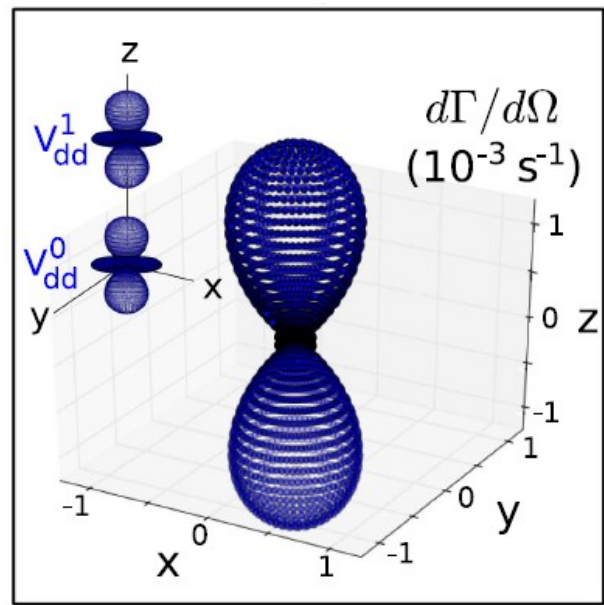
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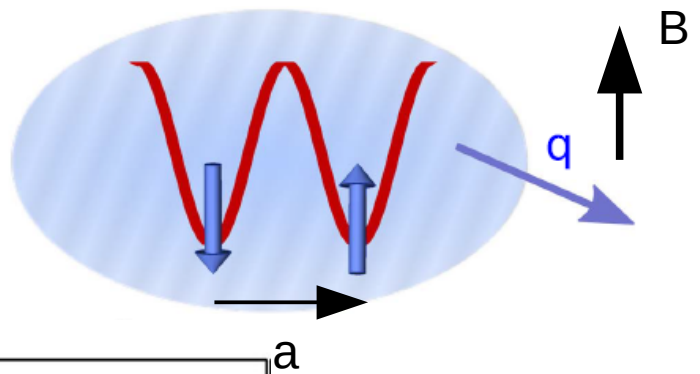


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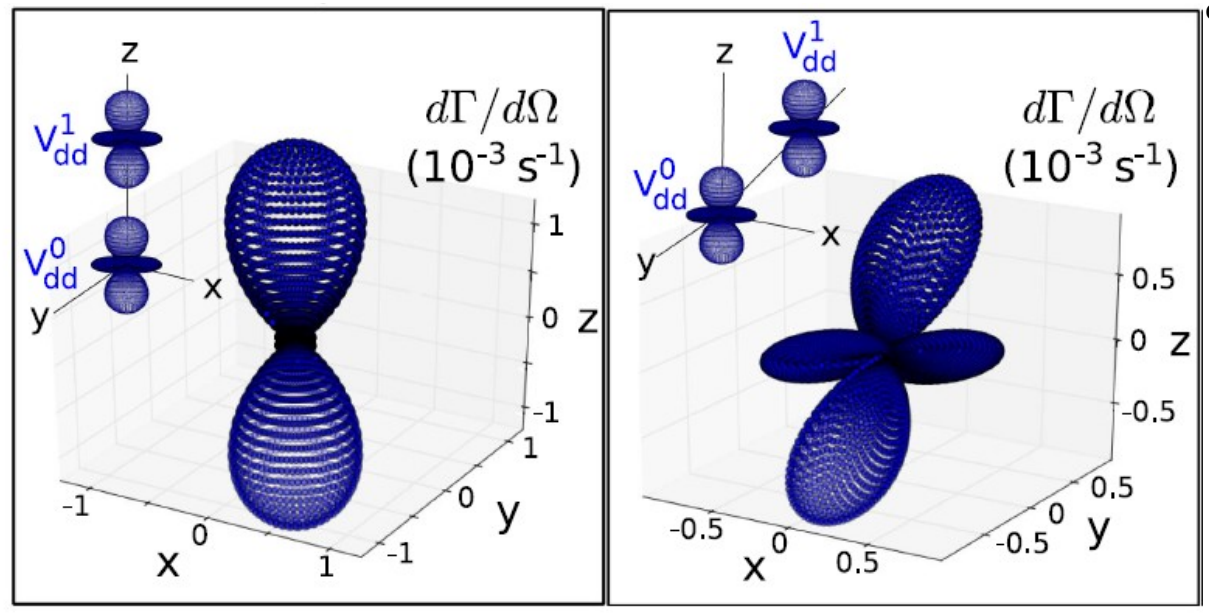
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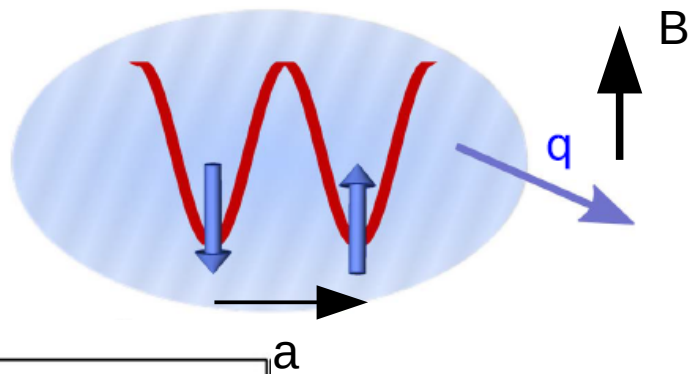


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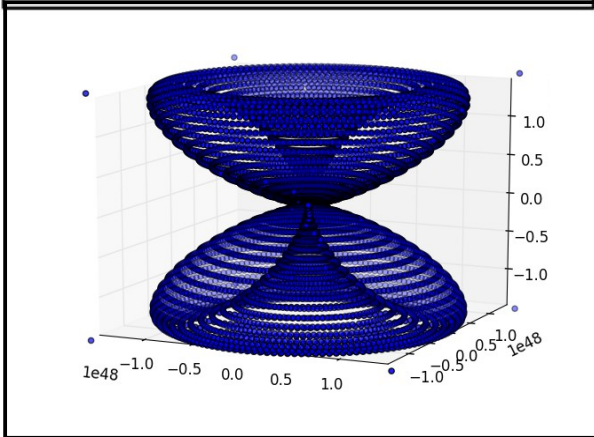
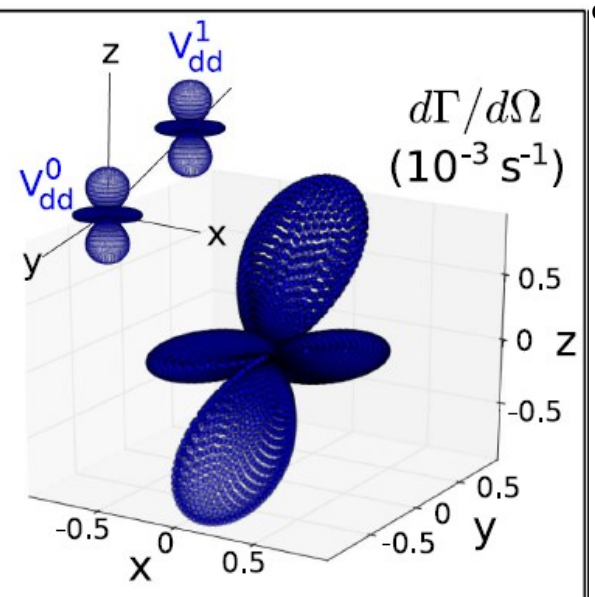
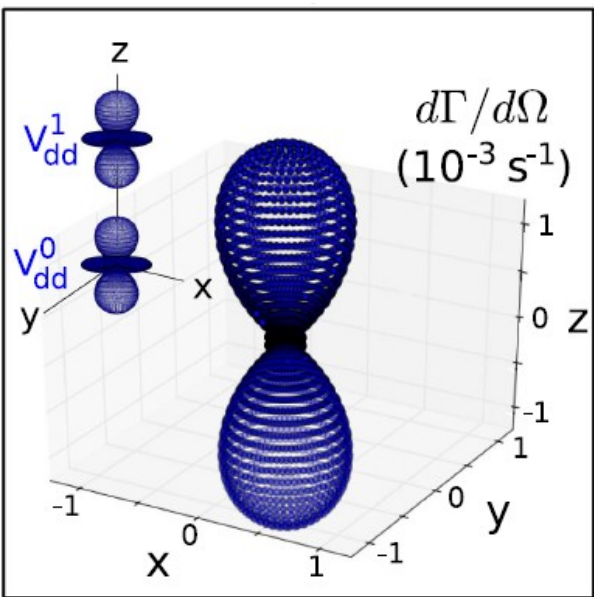
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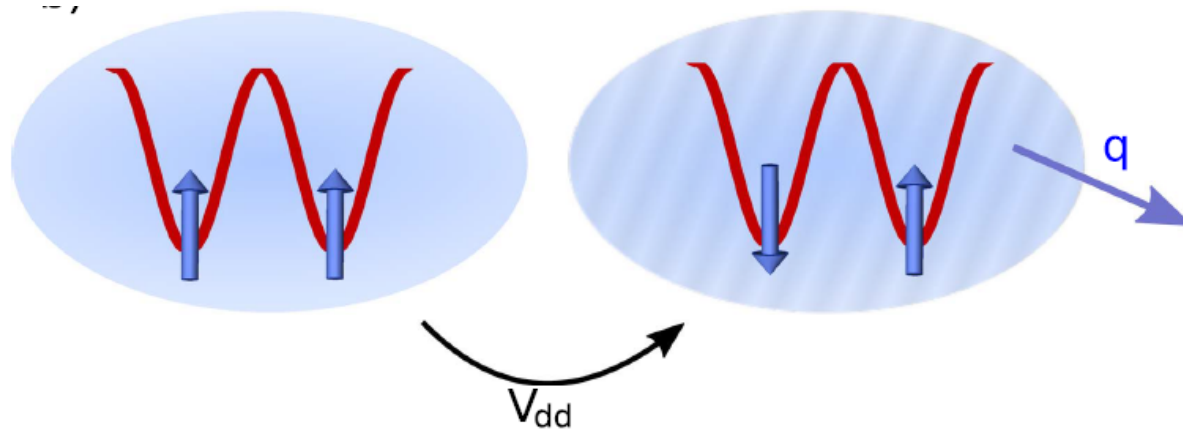
Severe anisotropy and collective spin dependence



from $S_{\text{tot}} = 1, m = 0$
to $S_{\text{tot}} = 0, m = 0$



from $S_{\text{tot}} = 1, m = \pm 1$
to $S_{\text{tot}} = 0, m = 0$



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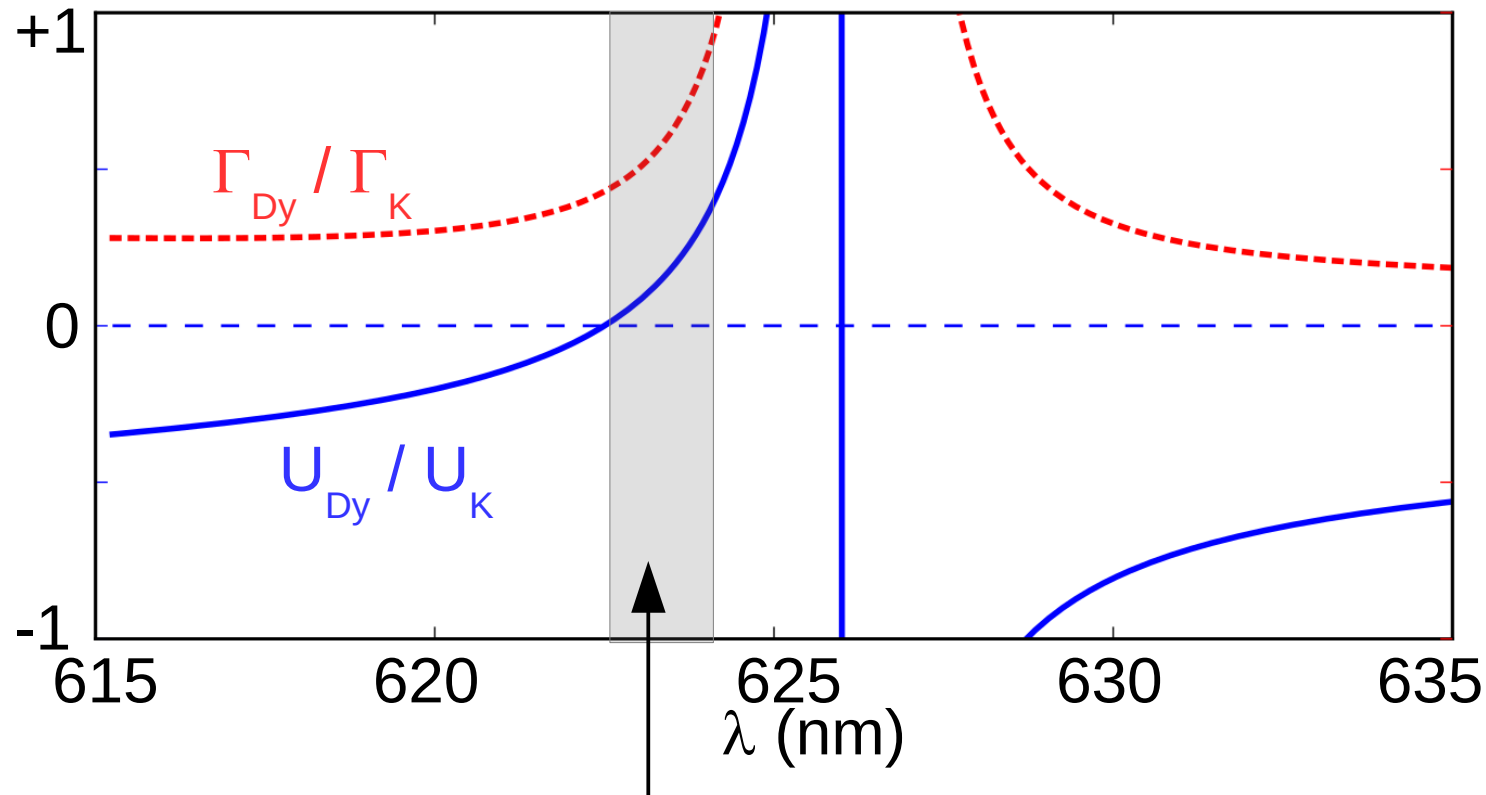
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Lattice potential: strong effect on the bath

⁴⁰K - ¹⁶⁴Dy

Given a lattice depth for the spin chain,
In the vicinity of 624 nm (Dy) **the lattice depth for the bath can be independently tuned**



- Mott regime for spin chain
- **3D coherence for bosonic bath**
- Light scattering sufficiently low

Anisotropy effects | **Dipolar bath stabilized from dynamical instabilities**
- **Enhanced interactions**

Convergence to a thermal state of the collective spin

^{40}K - ^{164}Dy

^{40}K , $F = 9/2$, restricted to $m = -9/2$ and $-7/2$, made degenerate ¹

$U_{\text{K}} = (25 \times 25 \times 3.5) E_r^{\text{K}}$ – **effective decoupled 1D chains**

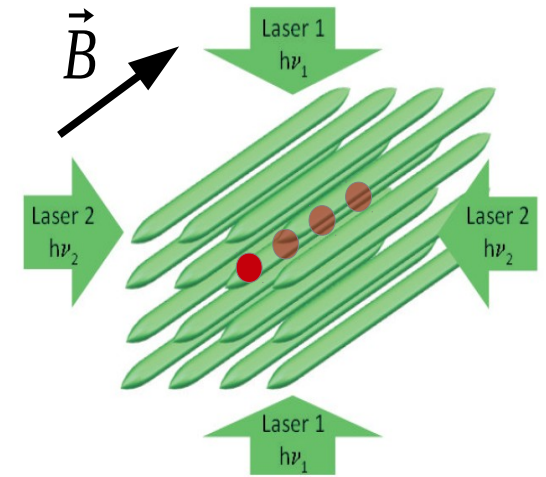
Weak axis : $U_{\text{int}}/t = 7.5$, $J = h \times 630 \text{ Hz} = k_B \times 30 \text{ nK}$

$U_{\text{Dy}} = (12 \times 12 \times 3.5) E_r^{\text{Dy}}$

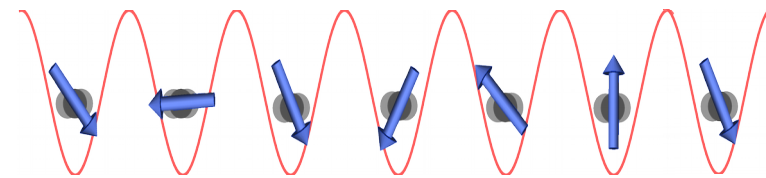
$\langle n_{\text{bec}} \rangle = 3.10^{13} / \text{cm}^3$

$T_{\text{BEC}} = 0,3 \text{ J} / k_B = 9 \text{ nK}$ [Trotzky 2010, Nat. Phys. **6**,998]

3D coherent BEC - Quantum depletion : 5 % [Xu 2006, PRL **96**, 180405]



Chain Length : 7



¹appendix ; and Gerbier 2006

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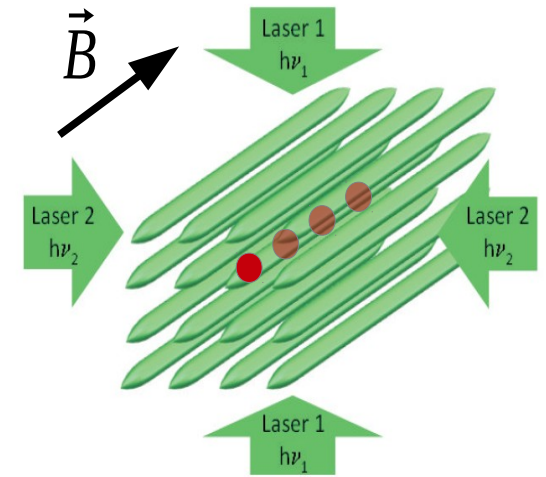
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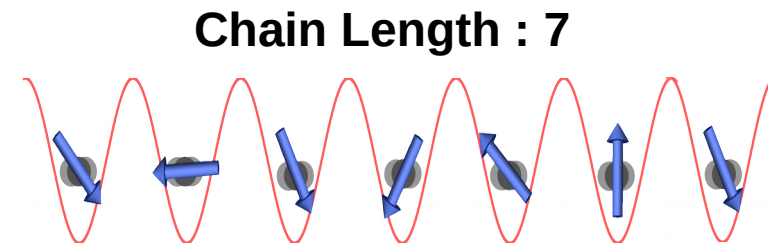
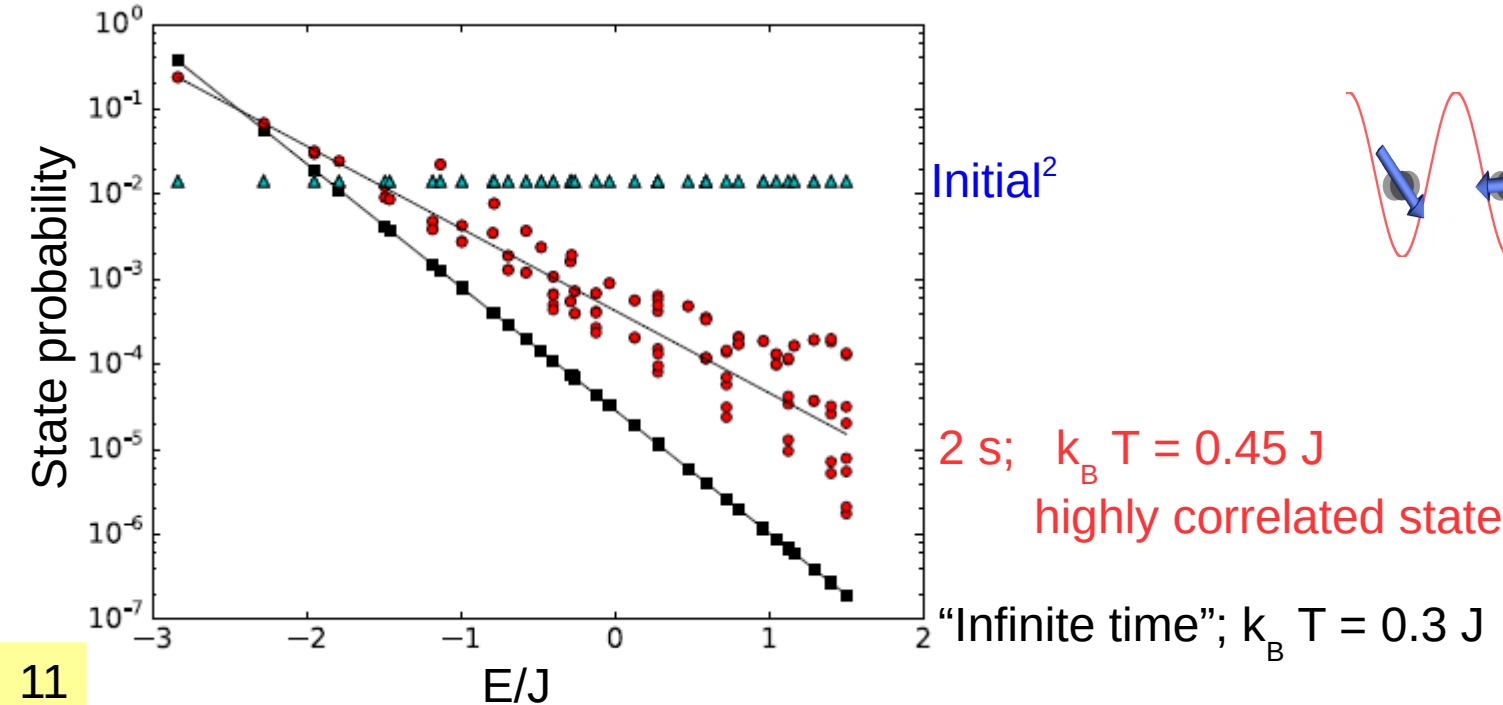
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Occupation of the $2^7 = 128$ spin chain eigenstates

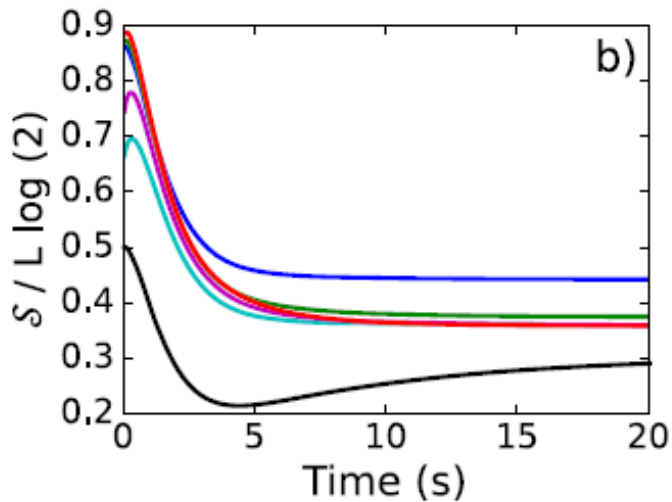
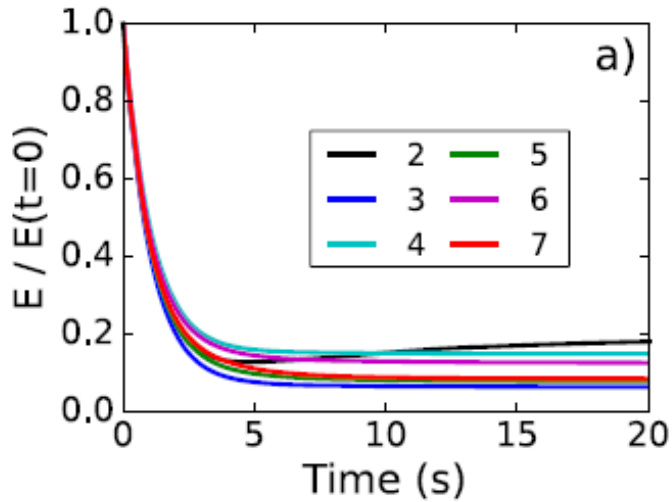


¹appendix ; and Gerbier 2006

²Here: initial magnetization 0

Collective spin dynamics

Initially balanced spin mixture



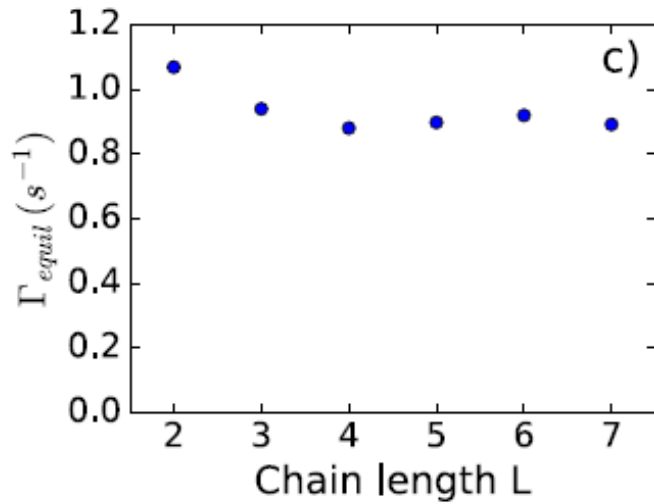
Energy

Von Neumann spin entropy

$$\mathcal{S} = -\sum_i p_i \log(p_i)$$

Collective spin dynamics

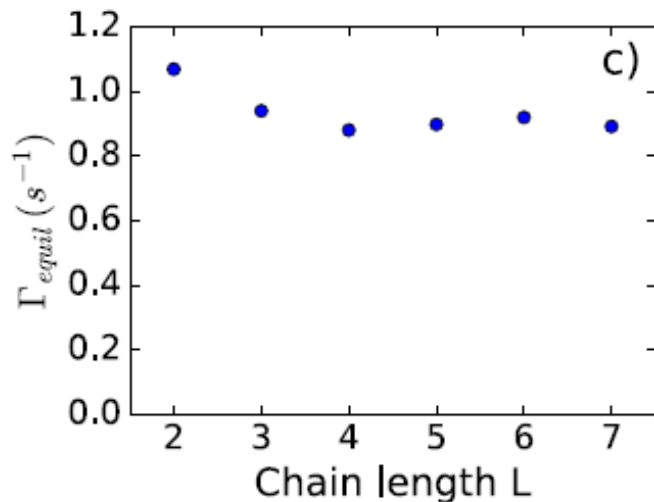
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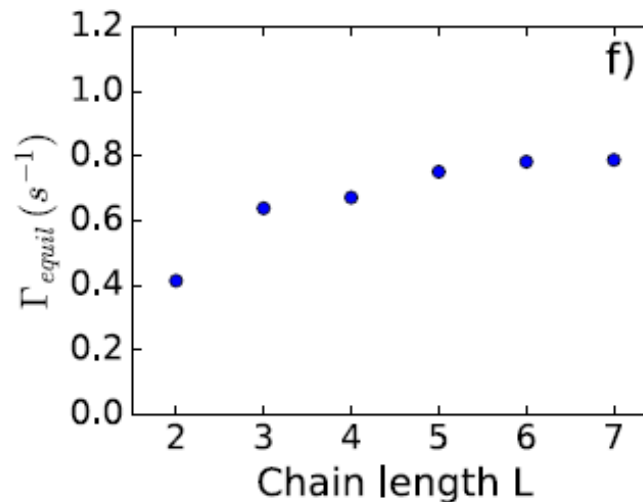
Equilibration rate tends to a value roughly independent on chain length and on preparation condition

Collective spin dynamics

Initially balanced spin mixture



Initially free spin mixture



Equilibration rate tends to a value roughly independent on chain length and on preparation condition

Timescale of order ~ 1 s – experimentally relevant, though not fast

Limited by restraining ourselves to very low quantum depletion (5%)

Faster dynamics plausible in deeper bath lattices,

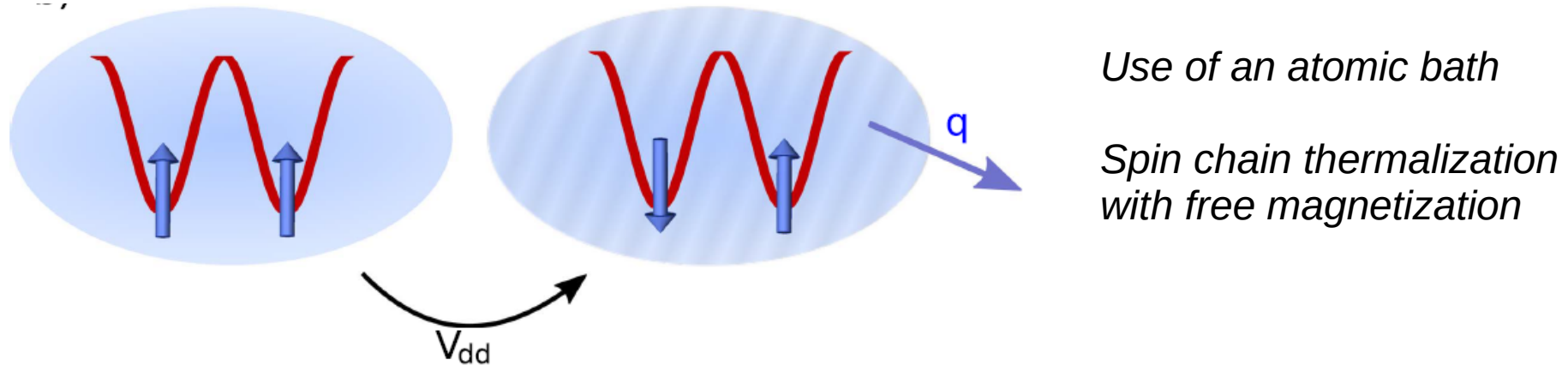
but this leaves the validity range of the Bogoliubov description

Dysprosium vs Erbium : about similar ($7\mu_B$, but also 583 nm lattice)

Alkali: ⁴⁰K has low Lande factor, but scientific interest of fermions for the t-J model

Conclusion and outlook

Dissipative preparation of strongly correlated spin states



The scheme relies on spin-orbit coupling in V_{dd}

→ **perspective**: cooling with a non-dipolar atomic bath using artificial SOC?

Spielman, Zwierlein, Zhang, Pan ...

A formalism describing dipole-coupled Mott spin chain and SF BEC in lattice

→ useful beyond Heisenberg chains (e.g., mixtures of dipolar isotopes in lattices)

Ferlaino, Lev, Pfau, Laburthe-Tolra, ...

→ other spinor species of interest (bosonic alkalis with higher Lande factor than ^{40}K)

→ **perspective**: set the formalism for a conducting **fermionic bath**
(large density of states at the Fermi energy, with large excitation wavevector)

Thank you for your attention

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Centre National de la Recherche Scientifique, Université Paris 13

ANR, DIM Nano'K, DIM Sirteq, IFRAF, IFCPAR

Magnetic Quantum gases group at LPL (Laburthe-Tolra's group):

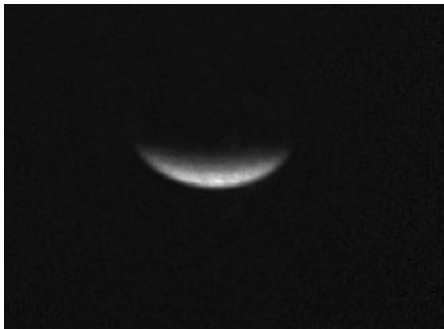
Two experiments : Strongly dipolar Chromium gases

SU($N \leq 10$) symmetric Strontium gases – new machine

One theory team on large spin quantum gases

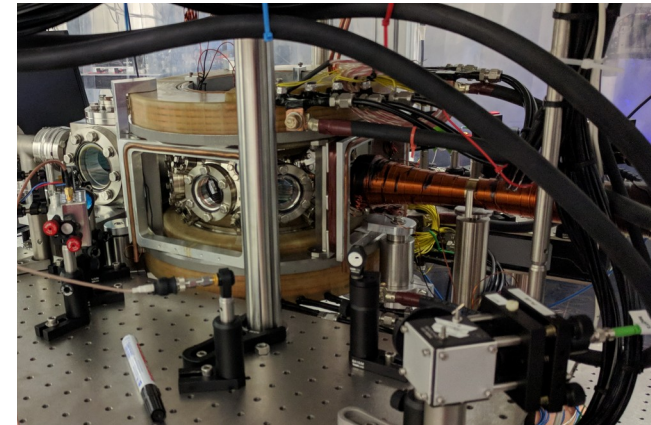
WE LOOK FOR A POST-DOC ON THE STRONTIUM MACHINE

A PHD ON THE CHROMIUM MACHINE

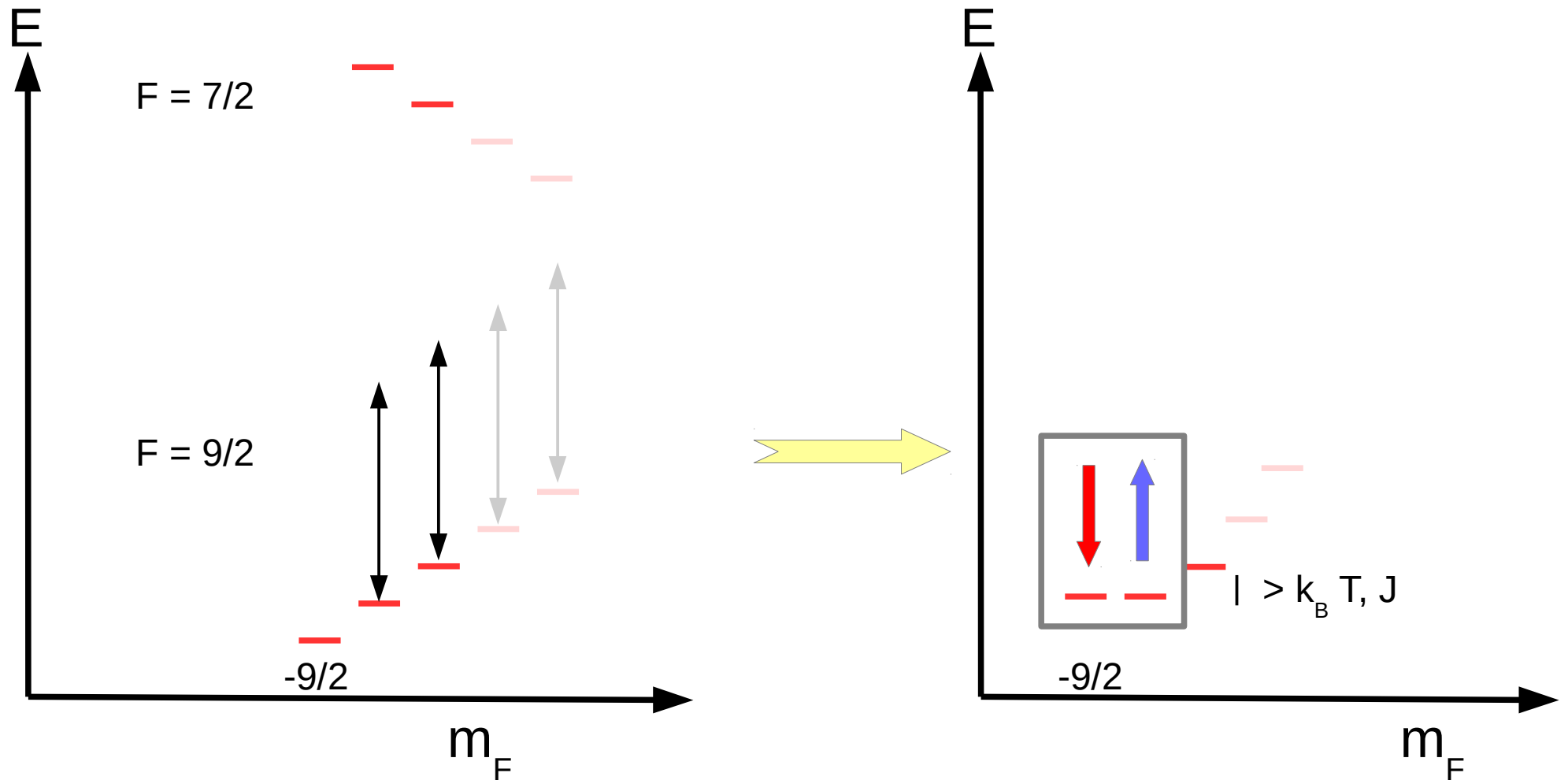


Narrow-line Sr MOT and
dipole trap: spring 2018

Objectives: SU(N) Quantum magnetism
Narrow-line manipulation tools



Microwave dressing



$B = 10$ mG : linear Zeeman effect with 3 kHz / m_F energy difference

Microwave dressing close to the hyperfine transition at 1,285 GHz, $F=9/2 \rightarrow F = 7/2$

Rabi frequency 100 kHz

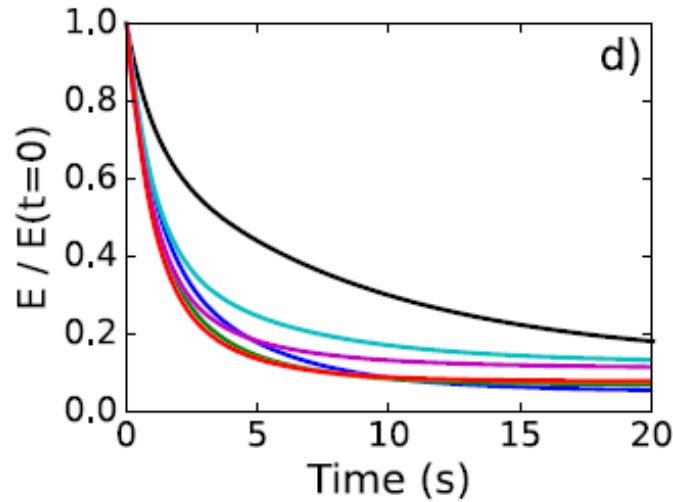
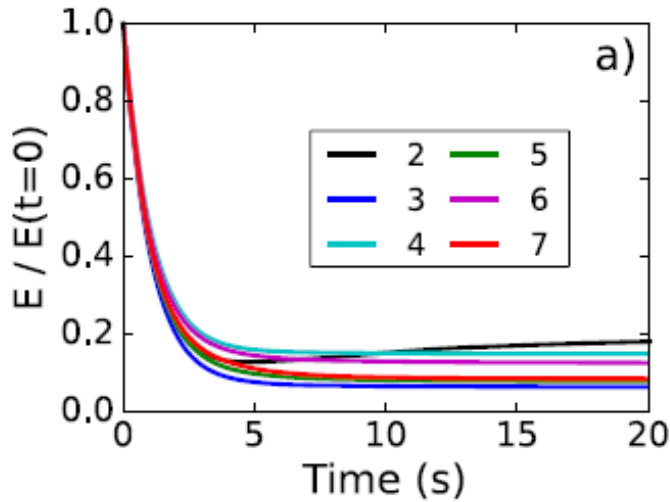
Detuning 830 kHz

Hyperfine state mixing 0,004

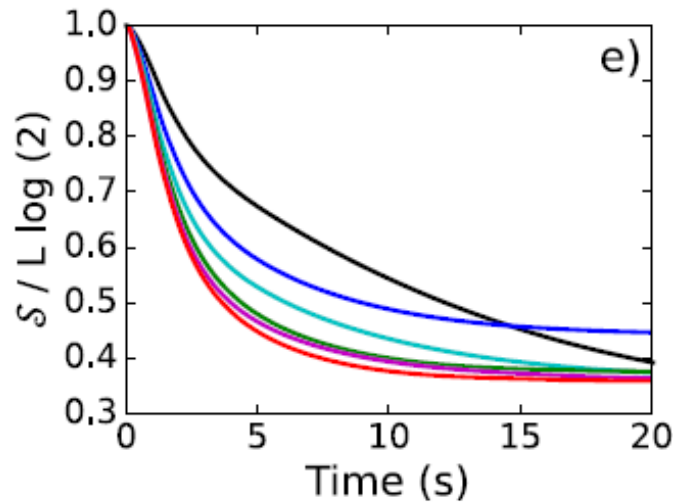
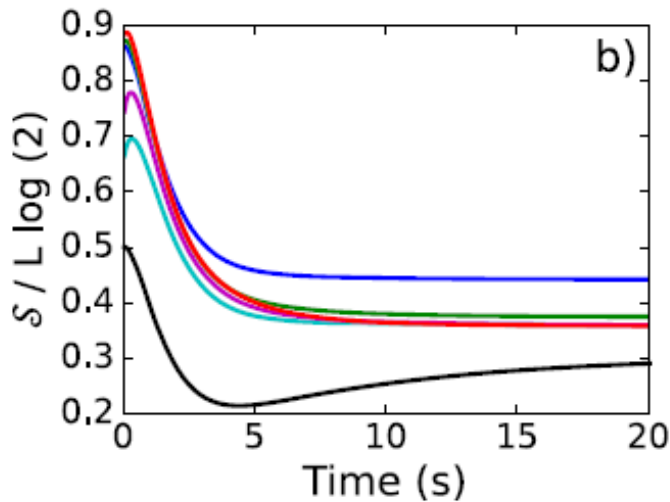
Collective spin dynamics

Initially balanced spin mixture

Initially free spin mixture



Energy



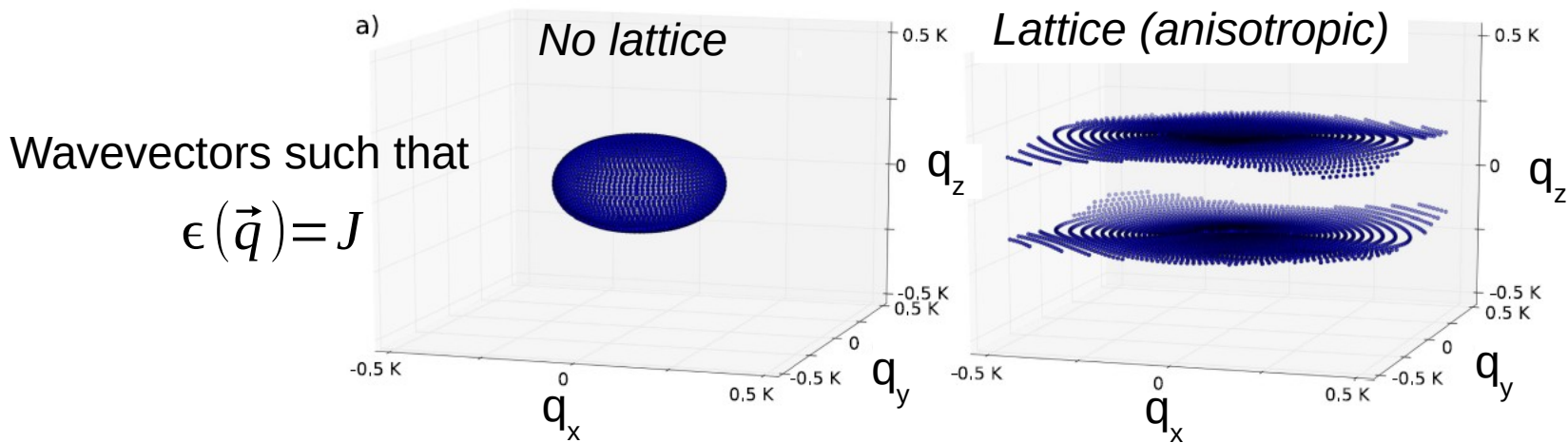
Von Neumann spin entropy

$$\mathcal{S} = -\sum_i p_i \log(p_i)$$

Lattice potential: strong effect on the bath

Enhanced interactions : very sensitive to anisotropies

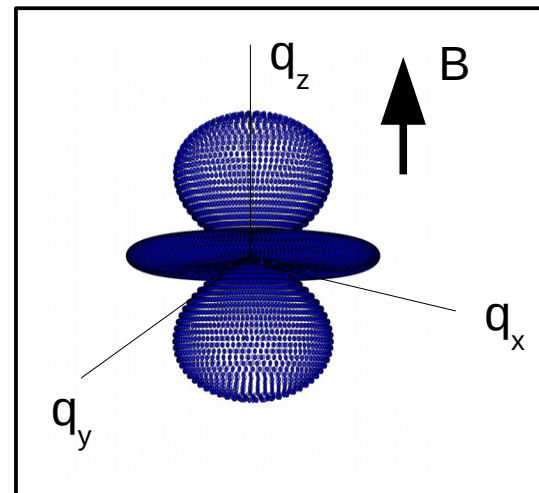
Dispersion relation : wavevectors and density of states



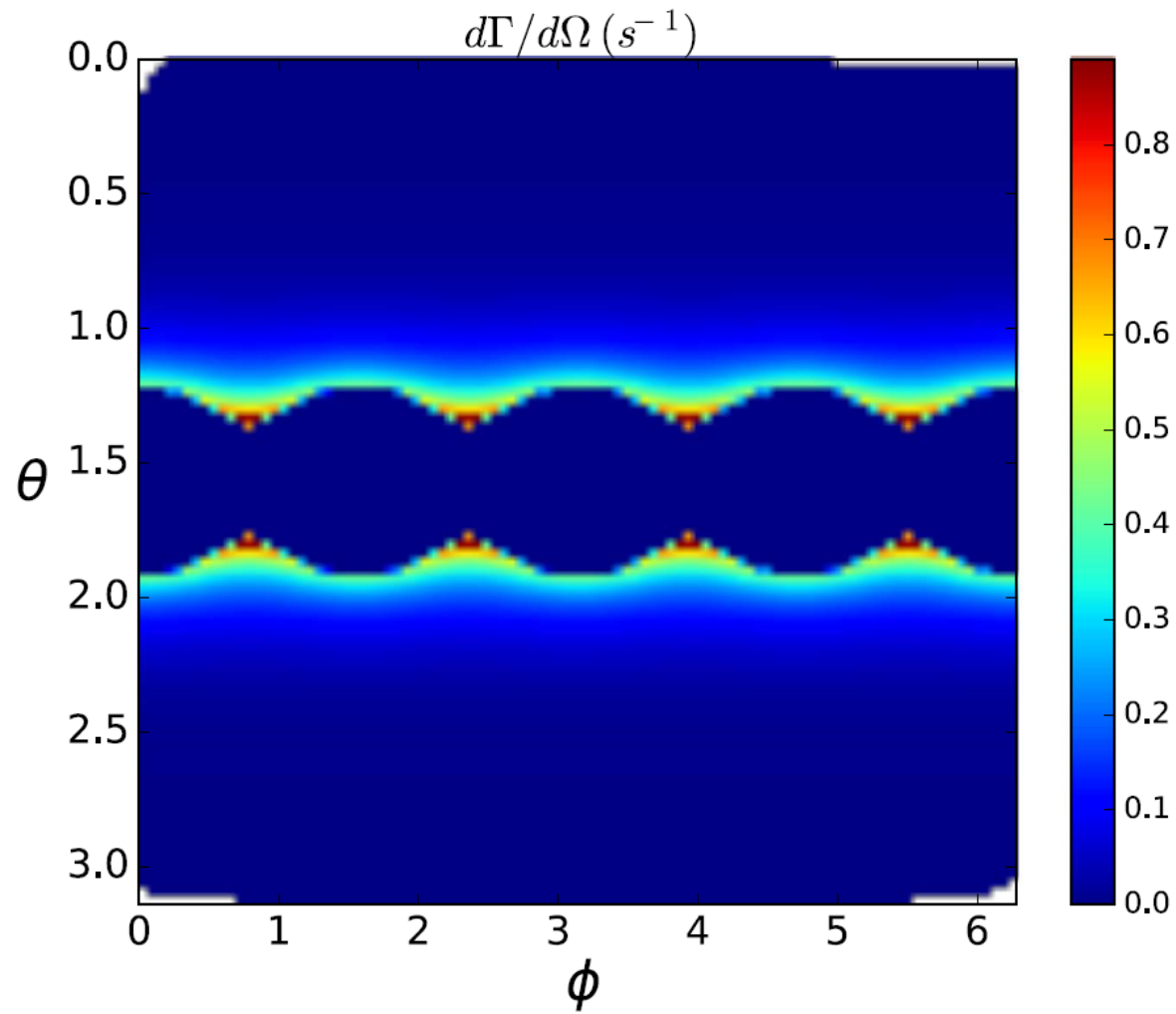
Coupling strength for a given mode \mathbf{q}

Mode decomposition onto plane waves vs Vdd anisotropy

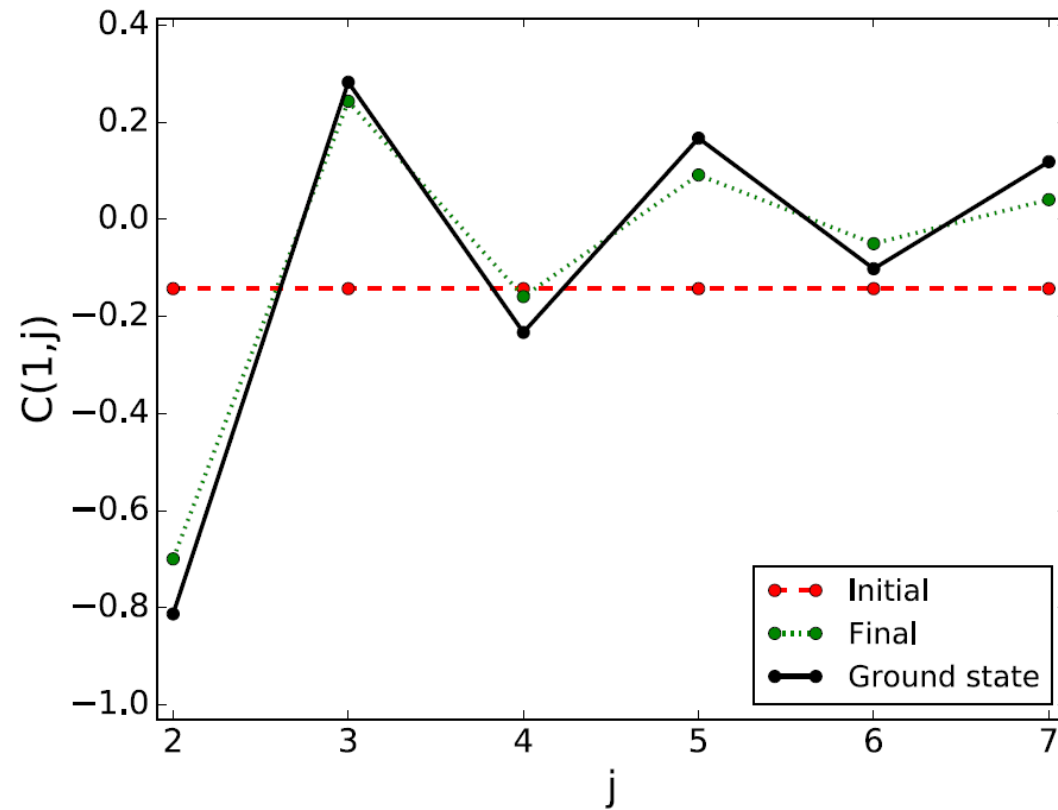
Vdd(q) :
Lobes of opposite sign



Radiation diagram in lattice

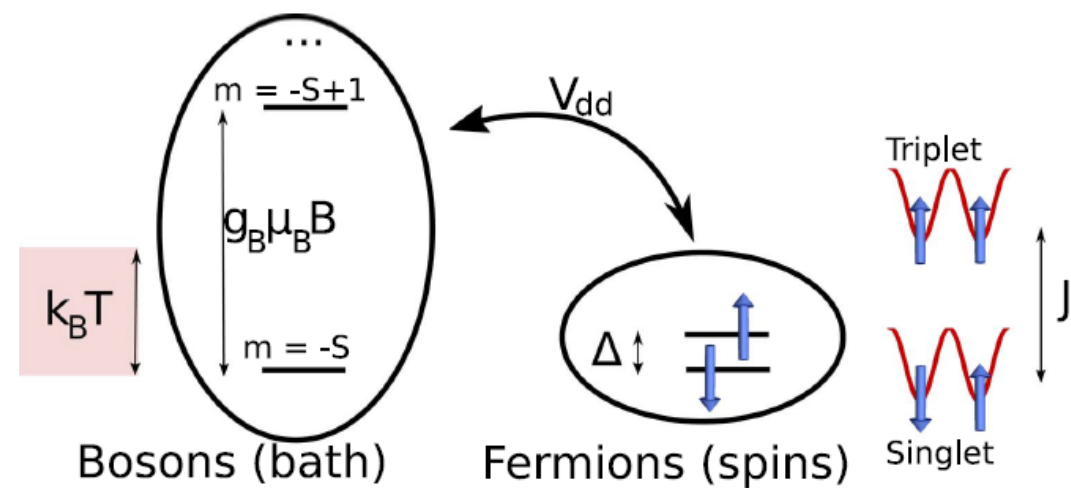


Correlations

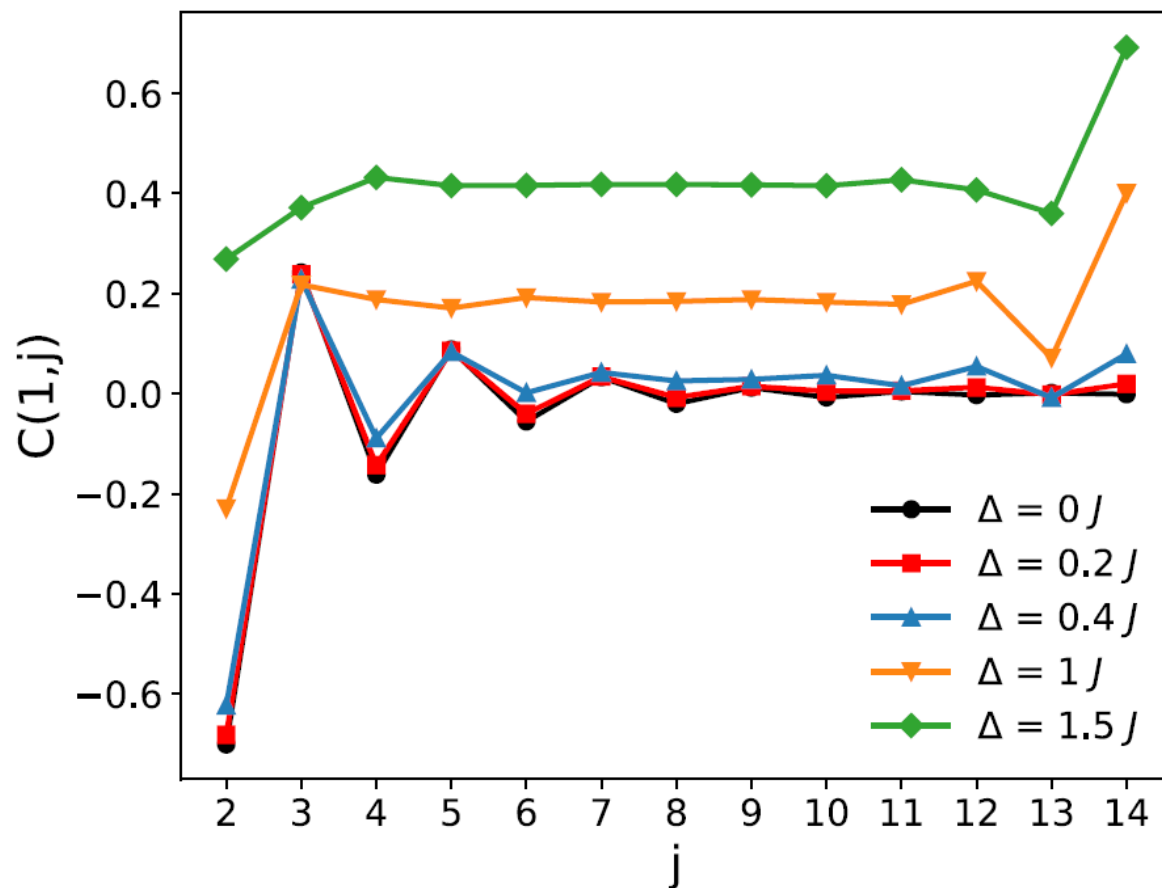


$$C(i, j) = 4 \langle \Sigma_{\text{eff},i}^z \Sigma_{\text{eff},j}^z \rangle$$

Robustness of the AF state to a bias Δ



$0.2 J \sim 0.4 \text{ mG}$



References

- M. Cazalilla, A. Ho, and T. Giamarchi, *New J. Phys.* **8**, 158 (2006)
- S. Diehl et. al., *Phys. Rev. Lett.* **105**, 227001 (2010)
- F. Gerbier et. al., *Phys. Rev. A.* **73**, 041602 (2006)
- Hart et. al., *Nature* **519**, 211 (2015)
- J. Kaczmarczyk, H. Weimer, and M. Lemeshko, *New J. Phys.* **18**, 093042 (2016)
- Mazurenko et al., *Nature* **545**, 462 (2017)
- Mathy et. al., *Phys. Rev. A* **86**, 023606 (2012).
- B. Pasquiou et. al., *Phys. Rev. A* **81**, 042716 (2010)
- B. Pasquiou et. al., *Phys. Rev. Lett.* **106**, 255303 (2011)
- A. Vogler et. al., *Phys. Rev. Lett.* **113**, 215301 (2014)